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Goodness of Generalized Seniority in Semi-magic Nuclei

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Abstract: Symmetry plays an important role in understanding the nuclear structure properties from the rotation of a nucleus to the spin, parity and isospin of nuclear states. This simplifies the complexity of the nuclear problems in one way or the other. Seniority is also a well known quantum number which arises due to the symmetry in the pairing interaction of nuclei. We present empirical as well as theoretical evidences based on decay rates which support the goodness of seniority at higher spins as well as in n -rich or, n -deficient nuclei. We find that the generalized seniority governs the identical trends of high-spin isomers in different semi-magic chains, where different set of nucleon orbitals from different valence spaces are involved.

Key words: generalized seniority; nuclear isomer; semi-magic nuclei

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1 Introduction

Understanding the complex nuclear structure of the atomic nucleus is an outstanding problem of nuclear physics. The pairing of nucleons is well known for several decades, which enables one in describing the ground states of even-even nuclei, and the extra stability of nuclei with an even number of nucleons than the nuclei with an odd number, *etc.* In atomic physics, pairing was understood in terms of seniority, first introduced by Racah in 1940s^[1]. In 1950s, Racah and Talmi^[2], and Flowers^[3] independently introduced the seniority scheme in nuclear physics. In simple terms, seniority may be defined as the number of unpaired nucleons in a given state, generally denoted as ν . It is now known that ν remains a good quantum number for states emanating from a pure- j configuration with $j \leq 7/2$. However, its validity for higher- j values has also been suggested.

Arima and Ichimura^[4], and Talmi^[5] further extended the seniority picture in single- j to the generalized seniority in multi- j . Generalized seniority takes care of the presence of multi- j orbitals in a given state. Semi-magic nuclei provide a fertile ground to study the various properties on the basis of the seniority scheme. One of the most interesting results from this scheme is the formation of seniority isomers. Seniority isomers are one of the well known categories of the nuclear isomers, *i.e.* the longer lived excited states, where the

hindrance to their decays have been explained in terms of seniority selection rules^[6]. It has been a general belief that the seniority isomers arise only in E2 decays between the same seniority states due to the vanishing decay probabilities at the mid-shell in seniority scheme. We have recently used the simple quasi-spin scheme to obtain the generalized seniority results in a mixed configuration coming from several degenerate orbitals, and applied it to the high-spin isomers in Sn isotopes^[7]. Hence, we have found for the first time odd-tensor E1 decaying 13^- isomers in Sn isotopes, a new category of isomers. We have then used the same scheme to understand the first excited 2^+ states in Sn isotopes and explained the asymmetric twin $B(E2)$ parabolas^[8].

In this paper, we have applied the “generalized seniority formalism” to the high-spin nuclear isomers in the semi-magic nuclei, particularly $Z = 50$ isotopes, $N = 82$ isotones and $Z = 82$ isotopes. We find that the generalized seniority remains a reasonably good quantum number for a set of states, particularly the high-spin isomers in these semi-magic nuclei. This further governs the identical behavior of these isomeric states in different semi-magic chains, having different set of active orbitals. We start by presenting a few empirical evidences and understanding the reasons behind, in section 2, which makes a good ground for the seniority calculations. Thereafter, we follow these empirical findings with our generalized seniority calculations and results for both even- and odd- A semi-magic

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nuclei. The overall conclusions of the present work have been presented in the last section.

2 Experimental evidences

The level schemes of the $^{119-130}\text{Sn}$ isotopes have been studied by using the reactions induced by light ions, deep inelastic reactions, or fission fragment studies by several researchers^[9-18]. Many isomer systematics have been identified for $N > 64$ Sn-isotopes, and the isomeric states 10^+ and $27/2^-$ have been characterized as seniority $v = 2$ and $v = 3$ states in these studies. Pietri *et al.*^[19] recently identified and confirmed the high-spin and high-seniority $v = 4$, 15^- isomeric state in ^{128}Sn . More recently, Astier *et al.*^[20, 21] reported detailed high-spin level schemes in the $^{119-126}\text{Sn}$ isotopes by using the binary fission fragmentation induced by heavy ions. Iskra *et al.*^[22] have also focused on high-seniority states in neutron-rich, even-even Sn-isotopes. It may be noted that there exists some deformed collective states giving rise to a full or a part of rotational band in the even-even light mass Sn isotopes with $A = 110 \sim 118$, interpreted as $2p - 2h$ proton configuration^[23-29]. But the 10^+ yrast isomeric states discussed in the present paper are not part of any rotational structure^[30]. More recently, the studies on Sn isotopes have been pushed much beyond the $N = 82$ shell closure and isomers in the $N = 86 \sim 88$ Sn-isotopes have been populated by Simpson *et al.*^[31] which shed a new light on the effective interaction in n-rich nuclei^[32].

The 10^+ and $27/2^-$ isomers have also been identified, in the $N = 82$ isotonic chain from $Z = 66$, Dy to $Z = 72$, Hf, as seniority $v = 2$ and $v = 3$ isomers coming from the $h_{11/2}$ proton orbital^[33]. Recently, the high-spin structure of five $N = 82$ isotones with $Z = 54 \sim 58$ has also been reported by Astier *et al.*^[34], where the 10^+ isomers have been described as broken pairs of protons from the $g_{7/2}$ and $d_{5/2}$ orbitals in the even-mass isotones.

We plot the excitation energies of the 10^+ isomers relative to 0^+ states and the $27/2^-$ isomers relative to $11/2^-$ states for the $Z = 50$ isotopes and the $N = 82$ isotones in the top and bottom panels of Fig. 1, respectively. It may be noted that the same valence orbitals are involved in both the $Z = 50$ and $N = 82$ chains. While the neutrons occupy these orbitals in the $Z = 50$ isomers, the protons take over the role in the $N = 82$ isomers. We find that all the main features observed in the $Z = 50$ isotopic chain are also present in the $N = 82$ isotonic chain and both appear to be nearly identical to each other. We can see that the energy gap is almost constant and particle num-

ber independent which is a well known signature of nearly good seniority^[35-37]. The 10^+ and the $27/2^-$ isomers, belonging to the even-even and even-odd nuclei respectively, are seen to follow each other very closely throughout the chains, if one puts the 0^+ and $11/2^-$ states on equal footing. This suggests that the nuclear configurations and structure for the 10^+ and the $27/2^-$ isomers should be very similar without any odd-even effect which also suggests the aligned nature of the involved nucleons.

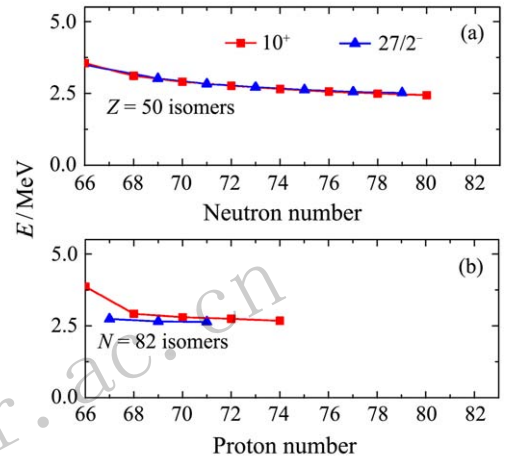


Fig. 1 (color online) Variation of the experimental energy values of the 10^+ and $27/2^-$ isomers in $Z = 50$ isotopes and $N = 82$ isotones.

We have plotted the measured half-lives (in μs) of these isomers with increasing nucleon numbers in the top and bottom panels of Fig. 2, for the $Z = 50$ and $N = 82$ chains, respectively. The half-lives of the 10^+ and $27/2^-$ isomers exhibit a rise near the middle of the active valence space (from neutron/proton numbers 66 to 82), attain a maximum value, and fall with increasing nucleon number. The 10^+ and $27/2^-$ isomeric states, for the $Z = 50$ isotopes, exhibit a maximum at the neutron numbers 72 and 73 respectively, where the $h_{11/2}$ neutron orbital becomes half-filled^[14, 17, 18]. On the other hand, for the $N = 82$ isomers, the peaks are observed at $Z = 70$ and 71 for the 10^+ and $27/2^-$ isomeric states respectively, where the $h_{11/2}$ proton orbital becomes half-filled^[33]. This happens because the electric quadrupole (E2) transition probabilities between a state J_i and another state J_f with same seniorities vanish at the middle^[7, 35, 36]. We can, therefore, foresee that the isomeric half-lives at the middle of the active valence space in the $Z = 50$ and $N = 82$ chains are most affected by the seniority selection rules. We also notice that the half-lives of odd- A Sn isotopes, *i.e.* the $27/2^-$ isomers are lower than the neighboring even- A Sn isotopes, *i.e.* the 10^+ isomers as expected from pairing consideration.

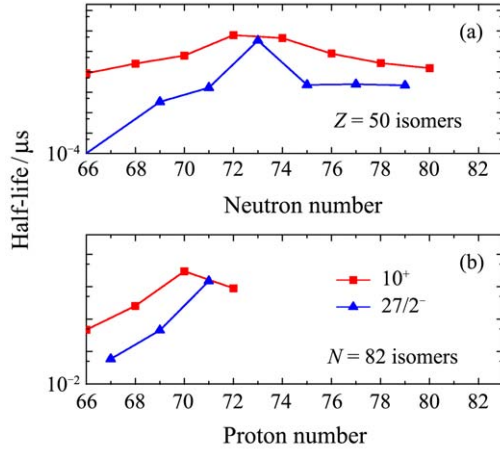


Fig. 2. (color online) Variation of the experimental half-life values of the 10^+ and $27/2^-$ isomers in $Z = 50$ isotopes and $N = 82$ isotones. The vertical scale is logarithmic.

We have also plotted the excitation energies of the 12^+ isomers relative to 0^+ states and the $33/2^+$ isomers relative to $13/2^+$ states in the top panel of Fig. 3 for the $Z = 82$, Pb isotopes. All the experimental data in this paper have been adopted from our atlas^[6], the ENSDF (Evaluated Nuclear Structure Data File)^[38], and the XUNDL (Unevaluated Nuclear Data List)^[39]

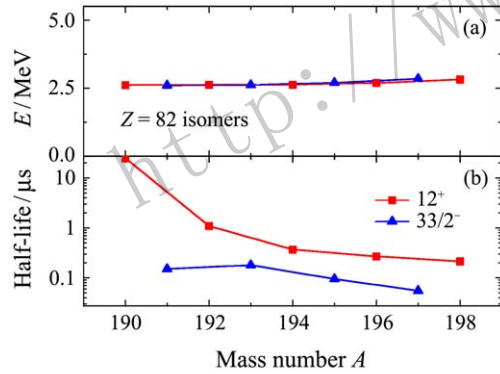


Fig. 3 (color online) Variation of the experimental energy and half-life values of the 12^+ and $33/2^+$ isomers in $Z = 82$ isotopes.

data sets. We can again see that both the Pb-isomers, even- A and odd- A , closely follow each other, and do not show any odd-even effect, as in the cases of Sn isotopes and $N = 82$ isotones. This again supports the empirical evidences of good seniority states with a particle number independent nature of the states. We, therefore, expect them to have similar origins in terms of their wave functions and nuclear configurations. We have also plotted their half-lives (in units of μs) in the bottom panel of Fig. 3, where one can observe that the half-lives show their increment towards mass number $A = 190$ ($N = 108$), which may be the middle of the active valence space for these Pb isomeric states.

We have further listed the experimental E2 gamma energies associated with the transitions $\Delta E_{0^+}^{2+}$ and $\Delta E_{11/2^-}^{15/2^-}$ for the even-even and odd- A Sn-isotopes in Table 1^[6, 20, 38-40]. The ratio of these transitions denoted as $R(15:2) = \Delta E_{11/2^-}^{15/2^-} / \Delta E_{0^+}^{2+}$ is observed to be ~ 1 for the $^{114-125}\text{Sn}$ isotopes. This suggests a complete alignment of the odd-neutron in the $h_{11/2}$ orbital, producing the $11/2^-$ spin state. This supports the observation that the $11/2^-$ state in odd- A Sn-isotopes and the 0^+ state in the neighboring even-even Sn-isotopes have great similarity in their wave functions. Similarly, the observed E2 gamma transitions $\Delta E_{10^+}^{12^+}$ and $\Delta E_{27/2^-}^{31/2^-}$ in even-even and odd- A Sn-isotopes have also been listed in Table 1 for $^{118-125}\text{Sn}$ -isotopes^[20, 40]. Fotiades *et al.*^[30] have compared the almost identical energies and similar structure involved in the $\Delta E_{0^+}^{2+}$ and $\Delta E_{10^+}^{12^+}$ γ -transitions within the same isotope for $^{116-126}\text{Sn}$, and suggested that the 10^+ isomeric state comes from the two aligned neutrons in the $h_{11/2}$ orbital. We have calculated the ratio of the transitions in odd- A Sn-isotope and its even-even core Sn-isotope, denoted as $R(31:12) = \Delta E_{27/2^-}^{31/2^-} / \Delta E_{10^+}^{12^+}$, which is also observed to be ~ 1 . The known gamma transition energies for the $N = 82$ isotones, have also been listed in Table 2^[40]. The ratios $R(15:2)$ and $R(31:12)$ again have the value ~ 1 , wherever these could be obtained.

Table 1 Comparison of the experimentally measured $\Delta E_{0^+}^{2+}$ and $\Delta E_{11/2^-}^{15/2^-}$ γ -transitions in even-even and odd- A Sn-isotopes for $N \geq 66$ and their ratio $R(15:2)$. Also, compared are the $\Delta E_{10^+}^{12^+}$ and $\Delta E_{27/2^-}^{31/2^-}$ γ -transitions involving states which decay to the 10^+ , $27/2^-$ isomers and their ratio $R(31:12)$. All the energies are in MeV.

Isotope	$\Delta E_{0^+}^{2+}$	$\Delta E_{10^+}^{12^+}$	Isotope	$\Delta E_{11/2^-}^{15/2^-}$	$\Delta E_{27/2^-}^{31/2^-}$	$R(15:2)$	$R(31:12)$
^{116}Sn	1.294		^{117}Sn	1.279		0.99	
^{118}Sn	1.230	1.237	^{119}Sn	1.220	1.179	0.99	0.95
^{120}Sn	1.171	1.190	^{121}Sn	1.151	1.083	0.98	0.91
^{122}Sn	1.141	1.103	^{123}Sn	1.107	1.043	0.97	0.95
^{124}Sn	1.132	1.047	^{125}Sn	1.088	0.924	0.96	0.88

Table 2 The same as Table 1, but for the $N = 82$ isotones with $Z \geq 66$. All the energies are in MeV.

Isotope	$\Delta E_{0^+}^{2^+}$	$\Delta E_{10^+}^{12^+}$	Isotope	$\Delta E_{11/2^-}^{15/2^-}$	$\Delta E_{27/2^-}^{31/2^-}$	$R(15:2)$	$R(31:12)$
^{148}Dy	1.677	1.932	^{149}Ho	1.560		0.93	
^{150}Er	1.578	1.446	^{151}Tm	1.478	1.332	0.94	0.92
^{152}Yb	1.531		^{153}Lu	1.432		0.94	

It is obvious from the observed values that the seniority v remains as 1 for the $11/2^-$ states, as coming from the unique parity $h_{11/2}$ orbital of the active valence space. We may also infer from these observations that the 10^+ and $27/2^-$ isomeric states, in both the chains, are maximally aligned decoupled states having similar wave functions and nuclear configurations. That is why the 10^+ and $27/2^-$ states closely follow each other in excitation energy without exhibiting any odd-even effect, as shown in Fig. 1. We, therefore, foresee the seniority as $v = 0$ for the 0^+ states, $v = 1$ for the $11/2^-$ states, $v = 2$ for the 2^+ states and $v = 3$ for the $15/2^-$ states. Similarly, we may assign the seniority $v = 2$ for the 10^+ states, $v = 3$ for the $27/2^-$ states, $v = 4$ for the 12^+ states, and $v = 5$ for the $31/2^-$ states for these n-rich Sn isotopes. The same seniority difference $\Delta v = 2$ between the $15/2^-$ and $11/2^-$ states, and

for the 2^+ and 0^+ states gives their corresponding ratio $R(15:2)$ as ~ 1 . The difference $\Delta v = 2$ also holds for the $31/2^-$ and $27/2^-$ states, and for the 12^+ and 10^+ states, which makes the ratio $R(31:12) \sim 1$.

We have also listed the experimental E2 gamma energies associated with the transitions $\Delta E_{0^+}^{2^+}$ and $\Delta E_{13/2^+}^{17/2^+}$ for the even-even and odd- A Pb-isotopes in Table 3^[6, 38, 39]. The ratio of these transitions denoted as $R(17:2) = \Delta E_{13/2^+}^{17/2^+} / \Delta E_{0^+}^{2^+}$ is observed to be ~ 1 for the $^{190-197}\text{Pb}$ isotopes. This suggests a complete alignment of the odd-neutron in the $i_{13/2}$ orbital, producing the $13/2^+$ spin state. This supports the observation that the $13/2^+$ state in odd- A Pb-isotopes and the 0^+ state in the neighboring even-even Pb-isotopes have great similarity in their wave functions, very similar to the isomers in other two semi-magic chains.

Table 3 The same as Table 1, but for the $Z = 82$ isotopes with $N \leq 108$. Comparison of the experimentally measured $\Delta E_{0^+}^{2^+}$ and $\Delta E_{13/2^+}^{17/2^+}$ γ -transitions in even-even and odd- A Pb-isotopes and their ratio $R(17:2)$. All the energies are in MeV.

Isotope	$\Delta E_{0^+}^{2^+}$	$\Delta E_{10^+}^{12^+}$	Isotope	$\Delta E_{13/2^+}^{17/2^+}$	$R(17:2)$
^{190}Pb	0.774	0.897	^{191}Pb	0.818	1.056
^{192}Pb	0.853	0.874	^{193}Pb	0.881	1.032
^{194}Pb	0.965	0.932	^{195}Pb	0.969	1.004
^{196}Pb	1.049	0.959	^{197}Pb	1.006	0.959

Similarly, the observed E2 gamma transitions $\Delta E_{12^+}^{14^+}$ have also been listed in Table 3 for the even-even $^{190-196}\text{Pb}$ isotopes; however, no measurements are available for $\Delta E_{33/2^+}^{37/2^+}$ in odd- A $^{190-197}\text{Pb}$ isotopes. We, therefore, could not calculate the ratio of the transitions in odd- A Sn-isotope and its even-even core Sn-isotope, denoted as $R(37:14) = \Delta E_{33/2^+}^{37/2^+} / \Delta E_{12^+}^{14^+}$, which is also expected to be ~ 1 . Therefore, similar arguments work for all the high-spin isomers in these $Z = 50$, $N = 82$ and $Z = 82$ semi-magic chains, while their respective valence spaces and intruder orbitals are different. This similarity hints towards the goodness of seniority. Keeping our previous results for even- A Sn isotopes in mind^[7], we have further done the generalized seniority calculations for decay rates to confirm the identical situation of various isomers from different nuclear regions.

3 Theoretical interpretation

We briefly present the formulas used in the $B(E2)$ calculations using the generalized seniority scheme, and successfully applied to the 10^+ , and 15^- isomers in Sn isotopes in^[7]. In this paper, we extend our studies to the 10^+ isomers in $N = 82$ isotones and the 12^+ isomers in $Z = 82$ isotopes. We further apply these results for the odd- A semi-magic nuclei, particularly for $27/2^-$, $19/2^+$, $23/2^+$ and $35/2^+$ isomers in Sn isotopes, $27/2^-$ isomers in $N = 82$ isotones and $33/2^+$ isomers in $Z = 82$ isotopes.

3.1 $B(E2)$ rates from generalized seniority

The $B(E2)$ values, between J_i and J_f states, in a mixed configuration $\tilde{j} = j \otimes j' \dots$ along with the corresponding total pair degeneracy $\Omega = \frac{1}{2}(2\tilde{j} + 1) = \frac{1}{2} \sum_j (2j + 1)$ by using the generalized seniority scheme

can be written as follows

$$B(E2) = \frac{1}{2J_i + 1} \times |\langle \tilde{j}^n v l J_f || \sum_i r_i^2 Y^2(\theta_i, \phi_i) || \tilde{j}^n v' l' J_i \rangle|^2 \quad (1)$$

This implies that the $B(E2)$ values show a parabolic behavior in the multi- j case depending upon the seniority of the states involved in the transition. We rewrite the seniority reduction formula for the reduced matrix elements with seniority conserving $\Delta v = 0$ transitions between the initial and final states for the completeness of the text in paper. The relations are as follows:

$$\langle \tilde{j}^n v l J_f || \sum_i r_i^2 Y^2(\theta_i, \phi_i) || \tilde{j}^n v' l' J_i \rangle = \left[\frac{\Omega - n}{\Omega - v} \right] \times \langle \tilde{j}^v v l J_f || \sum_i r_i^2 Y^2(\theta_i, \phi_i) || \tilde{j}^v v' l' J_i \rangle. \quad (2)$$

The $B(E2)$ values, which depends on the particle number n , the generalized seniority v and the corresponding total pair degeneracy Ω , can be calculated by using these formulas. These formulas take care of mixing of the active orbitals in the valence space. We present details of the calculations and results in the next section. Note that some information on radial integrals along with the matrix elements of spherical harmonics is hidden in the constant of proportionality.

3.2 Even- A Semi-magic nuclei

We have successfully shown that the high-spin E2 decaying 10^+ , 15^- and E1 decaying 13^- isomers are similar in their decay trends for even-even Sn isotopes, see Ref. [7] for details. We now apply the same formalism^[7] to the isomers in other semi-magic chains, particularly, the 10^+ isomers in $N = 82$ isotones, and the 12^+ isomers in $Z = 82$ isotopes. Both the isomeric chains decay by E2 transitions as shown in Fig. 5 and 6, and their measured $B(E2)$ trends are quite similar to the 10^+ isomers in the Sn-isotopes (follow Fig. 4). We choose the active valence space as $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ in the $N = 82$ isotones and $i_{13/2}$, $f_{7/2}$ and $p_{3/2}$ in the Pb-isotopes. The resultant \tilde{j} and Ω values become $17/2$ and 9 in the $N = 82$ isotones, similar as in the case of $Z = 50$ isotopes, since both the $Z = 50$ and $N = 82$ chains share same orbitals in their active valence space of particle number $50 \sim 82$. We consider $Z = 64$ as core by assuming the $g_{7/2}$ and $d_{5/2}$ orbitals as completely filled.

Therefore, we fit the $n = 2$ situation at $Z = 66$, ^{146}Dy for the seniority $v = 2$ isomers in the $N = 82$ chain, and get the resultant parabolic trend as shown in Fig. 5. The calculated trends reproduce the experimental data quite well; $\Omega = 6$ gives the best fit to the data, a different situation as compared to the 10^+ ,

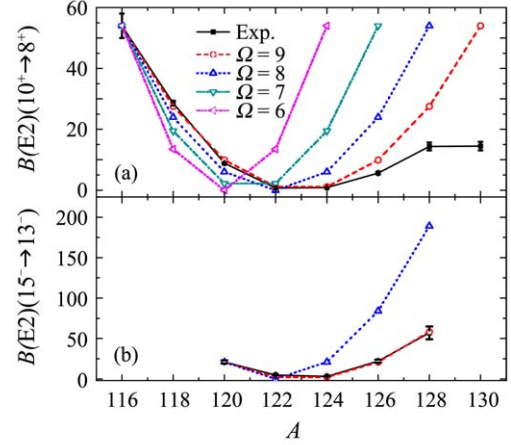


Fig. 4 (color online) Variation of the $B(E2)$ values of the 10^+ and 15^- isomers in $Z = 50$ isotopes^[7]. All the values are shown in the units of $e^2 \text{ fm}^4$.

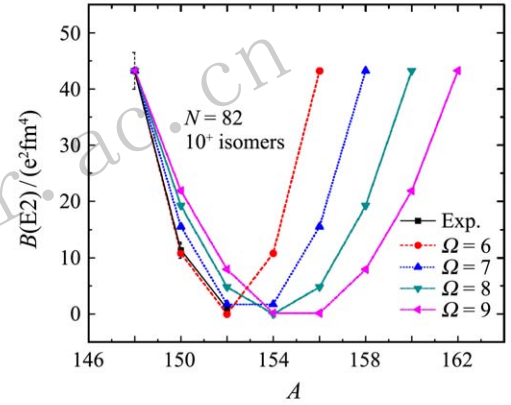


Fig. 5 (color online) Variation of the $B(E2)$ values of the 10^+ isomers in $N = 82$ isotones.

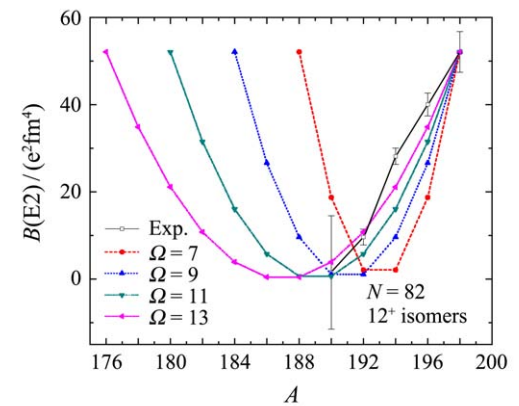


Fig. 6 (color online) Variation of the $B(E2)$ values of the 12^+ isomers in Pb-isotopes.

$Z = 50$ isomers, where $\Omega = 9$ gives the best fit. This can be understood in terms of different valence particles, particularly the involvement of protons in the case of $N = 82$ isotones. It appears to be closer to the pure seniority scheme, as quoted in the previous literature^[35].

On the other hand, the resultant \tilde{j} and Ω values become 25/2 and 13 for the Pb-isotopes, where $\tilde{j} = i_{13/2} \otimes f_{7/2} \otimes p_{3/2}$. The active valence space of $N = 82 \sim 126$ for these isotopes consists of $h_{9/2}$, $i_{13/2}$, $f_{7/2}$, $p_{3/2}$, $f_{5/2}$ and $p_{1/2}$ orbitals. We consider $h_{9/2}$ as completely filled, so the next three active orbitals $i_{13/2}$, $f_{7/2}$ and $p_{3/2}$ will be filled at ^{200}Pb . We fix the proportionality constants by fitting the measured values for ^{198}Pb for the seniority $v = 2$ isomers in the $Z = 82$ chains, respectively. The calculated results reproduce the experimental trends for both these isomeric chains as shown in Fig. 6. The calculations have been done by considering the transitions as seniority conserving ones ($\Delta v = 0$) along with generalized seniority of these states as $v = 2$. One can also expect the occurrence of these generalized seniority $v = 2$ isomers in the gaps of experimental data for both the $N = 82$ and $Z = 82$ chains extending towards proton-deficient and neutron-deficient sides, respectively (follow the Figs. 5 and 6 for the same).

One can, therefore, observe the identical behavior of the $B(E2)$ s in the high-spin isomers for all the three semi-magic chains, $Z = 50$, $N = 82$, and $Z = 82$. This is due to their identical generalized seniorities and the $\Delta v = 0$ transitions. It is interesting to note that the valence particles in $Z = 50$ chain are neutrons while protons become active in the case of $N = 82$ chain, and they share same orbitals in the active valence space of $50 \sim 82$. On the other hand, the $Z = 82$ chain has a different neutron valence space along with different set of active orbitals. In spite of these differences with each other, high-spin isomers in all the three chains appear to follow the same microscopic scheme of generalized seniority. This highlights the importance of configuration mixing required in the generation of these states. Hence, the generalized seniority behaves almost as a good quantum number for these states in all three semi-magic chains.

We note that the seniority $v = 4$ isomers are only known in the $Z = 50$ isotopes up to now. Due to the strong validity of generalized seniority in all the three semi-magic chains, one can expect and predict the high-spin and high-seniority $v = 4$ isomers in the remaining two chains as well. Measurements in this direction should be made to confirm this scenario.

3.3 Odd- A Semi-magic nuclei

We now study the $27/2^-$, $19/2^+$ and $23/2^+$ isomers of odd- A Sn-isotopes in light of the generalized seniority scheme. We assume that the $g_{7/2}$ and $d_{5/2}$ orbitals are completely filled up to ^{114}Sn ; hence, the remaining active orbitals for the Sn-isomers are $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals in the 50-82 valence space. We

have performed the generalized seniority calculations assuming $v = 3$, and $\Delta v = 0$ transitions and by fitting the value of ^{119}Sn isotope, as shown in Fig. 7. We calculate the $B(E2)$ values for the $27/2^-$ isomers using Ω values of 7, 8 and 9 corresponding to $\tilde{j} = h_{11/2} \otimes s_{1/2}^2$, $\Omega = 7$; $\tilde{j} = h_{11/2} \otimes d_{3/2}^2$, $\Omega = 8$; $\tilde{j} = h_{11/2} \otimes d_{3/2} \otimes s_{1/2}$, $\Omega = 9$, respectively. We present the calculated and experimental $B(E2)$ values for the $27/2^-$ isomers in Fig. 7 for comparison. One can see that the calculated values for $\Omega = 9$ fit the experimental data reasonably well. This confirms that the $27/2^-$ isomers behave as generalized seniority $v = 3$ isomers having $\Delta v = 0$ transitions to the lower lying $23/2^-$ states, and support the mixing of all the three $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals.

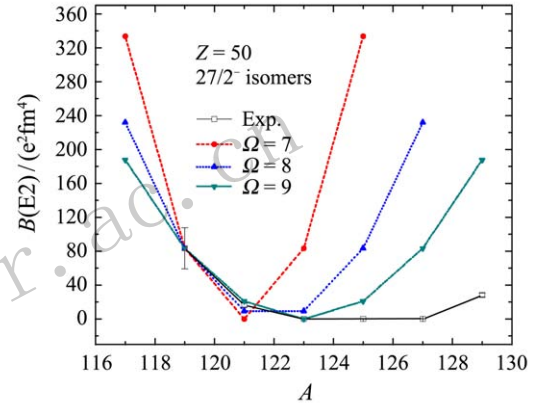


Fig. 7 (color online) Variation of the $B(E2)$ values of the $27/2^-$ isomers in Sn-isotopes^[42].

We also plot the $B(E2)$ values vs. A (mass number) for the $19/2^+$ and $23/2^+$ isomers, in the top and bottom panels of Fig. 8, respectively. Experimental data for these isomers have been taken from the recent measurements of Iskra et al.^[41] and the references therein. We find that the calculated results from $\Omega = 9$ value, and generalized seniority $v = 3$ with $\Delta v = 0$ transitions, are again able to explain the experimental trend reasonably well. Note that we fit the experimen-

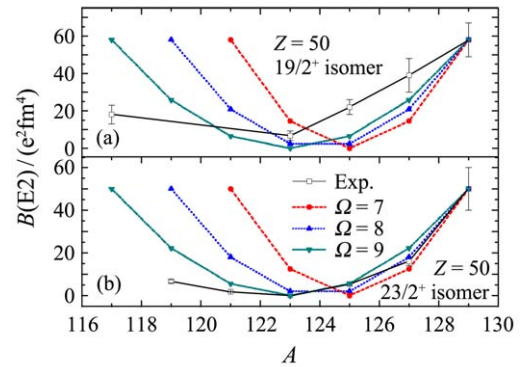


Fig. 8 (color online) Variation of the $B(E2)$ values of the $19/2^+$, and $23/2^+$ isomers in Sn-isotopes^[42].

tal value of ^{129}Sn in calculations for both the $19/2^+$ and $23/2^+$ isomers. This confirms that the $19/2^+$, $23/2^+$ isomers decay to the same seniority states, and are generalized seniority $v = 3$ isomers (involving $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals). Note that the seniority $v = 4$, 13^- , 15^- isomers, and seniority $v = 2$, 10^+ isomers in the even-even $^{116-130}\text{Sn}$ isotopes have already been explained by using the same mixed configuration having Ω value of 9^[7]. We can, therefore, conclude that the high-spin isomers in $^{116-130}\text{Sn}$ mass region arise from the mixing of all the available $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals for both even and odd- A isotopes.

On the other hand, we find that the inclusion of $d_{5/2}$ orbital along with $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals is required for explaining the first 2^+ states in these Sn-isotopes^[8]. However, the configuration changes for the first excited 3^- states showing the octupole character having $d-h$ orbitals^[42]. It is quite obvious that the dominance of $h_{11/2}$ orbital increases, while going towards the high-spin and high-seniority states. This type of information by generalized seniority guides us to infer the nature of configuration mixing and wave functions involved in the generation of a given set of states. Also, the structure information for both even- and odd- A nuclei are similar, irrespective of the mass region, involved valence particles, orbitals, *etc.*, except for an extra nucleon in odd- A systems. To sum up, we find that the seniority and generalized seniority provides a finger-print evidence for the similarity between various isomers in the semi-magic chains.

On the basis of this interpretation, we have also analyzed the single measured value at ^{123}Sn , for the higher seniority $v = 5$, ($35/2^+$) isomer. We have fitted the value of ^{123}Sn , and calculated the values for the other neighboring isotopes assuming $\Delta v = 0$ transitions using $\Omega = 8$ and the corresponding mixed configuration. We have taken ^{114}Sn as core; this implies that the first location ($n = 1$) to have seniority $v = 5$ state corresponds to ^{119}Sn . The $\Omega = 9$ and the related mixed configuration can not be fitted using the experimental value at ^{123}Sn as it leads to a zero value for the coefficient $(\frac{\Omega-n}{\Omega-v})^2$ in the middle ($n = v = 5$). We have plotted in Fig. 9, the $B(E2)$ trend for the seniority $v = 5$, $35/2^+$ isomer using $\Omega = 8$ value and the respective mixed configuration. These calculations, therefore, help us to predict some unknown values also. It is quite obvious that these high-seniority $v = 5$, $35/2^+$ isomers in odd- A n-rich Sn-isotopes can be related to the seniority $v = 4$ isomers in even- A n-rich Sn-isotopes.

We further study the $27/2^-$ isomers in odd- A $N = 82$ isotones using the same generalized seniority scheme. We have already pointed out that the Sn isotopes and $N = 82$ isotones have same set of orbitals

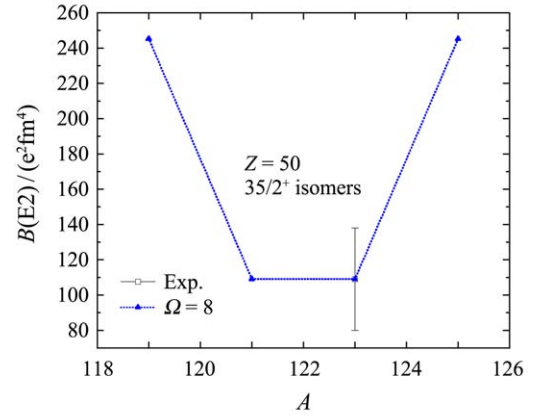


Fig. 9 (color online) Variation of the $B(E2)$ values of the $35/2^+$ isomers in Sn-isotopes^[42].

in the active valence space of nucleon number $50 \sim 82$, while neutrons are active in Sn isotopes and protons are active in the $N = 82$ isotonic chain. We assume that the $g_{7/2}$ and $d_{5/2}$ orbitals are completely filled up to ^{146}Gd ($Z = 64$); hence, the remaining active orbitals for these $N = 82$ isomers are $h_{11/2}$, $d_{3/2}$ and $s_{1/2}$ orbitals. We have performed the generalized seniority calculations assuming $v = 3$, and $\Delta v = 0$ transitions and by fitting the value of ^{149}Ho isotone, as shown in Fig. 10. We calculate the $B(E2)$ values for the $27/2^-$ isomers using Ω values of 7, 8 and 9 corresponding to $\tilde{j} = h_{11/2} \otimes s_{1/2}$, $\Omega = 7$; $\tilde{j} = h_{11/2} \otimes d_{3/2}$, $\Omega = 8$; $\tilde{j} = h_{11/2} \otimes d_{3/2} \otimes s_{1/2}$, $\Omega = 9$, respectively, as in the case of Sn isotopes. We present the calculated and experimental $B(E2)$ values for the $27/2^-$ isomers in Fig. 10 for comparison. One can see that the calculated values for $\Omega = 6$ best fit the experimental data. This confirms that the $27/2^-$ isomers behave as more like pure seniority $v = 3$ isomers having $\Delta v = 0$ transitions to the lower lying $23/2^-$ states, and support the dominance of $h_{11/2}$ orbital. The same argument has already been shown to work for even- A $N = 82$ isotones too.

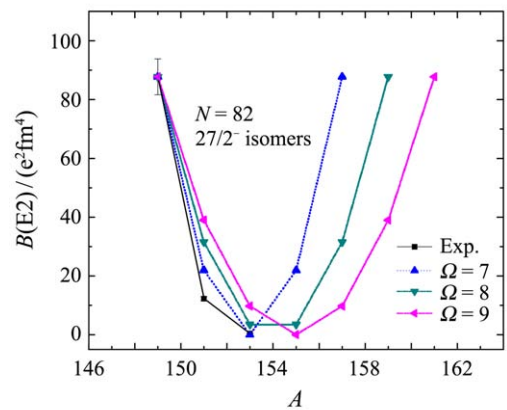


Fig. 10 (color online) Variation of the $B(E2)$ values of the 10^+ isomers in $N = 82$ isotones.

We next present a comparison of the $33/2^+$ isomers in Pb-isotopes on the same footing, where we again calculate the $B(E2)$ values for these isomers using $i_{13/2}$, $f_{7/2}$ and $p_{3/2}$ as active orbitals. We use the possible mixed configurations: $\Omega = 7$ ($i_{13/2}$), $\Omega = 9$ ($i_{13/2} \otimes p_{3/2}^2$), $\Omega = 11$ ($i_{13/2} \otimes f_{7/2}^2$), and $\Omega = 13$ ($i_{13/2} \otimes p_{3/2} \otimes f_{7/2}$). We assume that the lowest lying $h_{9/2}$ orbital is full, and therefore, the other three orbitals ($i_{13/2} \otimes p_{3/2} \otimes f_{7/2}$) become full at ^{200}Pb . Since we do not have experimental data at lower mass neutron-deficient side, we fit the value of ^{197}Pb by assuming ^{200}Pb as the completely full configuration, using generalized seniority $v = 3$ for $\Delta v = 0$ transitions. The calculated values from $\Omega = 13$ having mixing of all the three active orbitals are able to explain the experimental data quite closely (See Fig. 11).

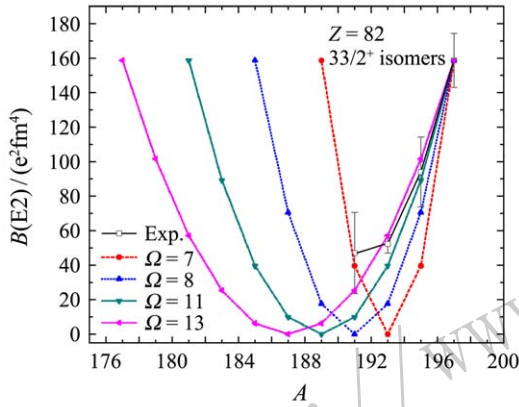


Fig. 11 (color online) Variation of the $B(E2)$ values of the $33/2^+$ isomers in Pb-isotopes^[42].

We, therefore, conclude that the $33/2^+$ isomers are seniority $v = 3$ isomers in Pb-isotopes, similar to the $27/2^-$ isomers in Sn-isotopes. Their similar $B(E2)$ trends can be explained in terms of the involvement of same seniority and generalized seniority, though coming from the different orbitals and different valence spaces. However, more measurements are required to obtain experimental data for the remaining nuclei and complete the picture. The same argument has been presented for the isomers in even- A nuclei. The seniority and generalized seniority hence play a unifying role in explaining the similar behavior in different sets of semi-magic nuclei.

4 Conclusion

We have used the quasi-spin formalism for degenerate multi- j orbitals to calculate the reduced electric transition probabilities in the semi-magic isomers. We find that the configuration mixing is essential to fully describe the $v = 2$, 10^+ isomers in the even- A Sn-isotopes, and the $v = 3$, $27/2^-$, $19/2^+$, $23/2^+$ and

$v = 5$, $35/2^+$ isomers in the odd- A Sn-isotopes and to explain the $B(E2)$ values in all these Sn-isomers. This formalism reproduces the experimental trend quite well and is also capable to predict some numbers for the gaps in the measurements. On the other hand, the situation for the $v = 2$, 10^+ and the $v = 3$, $27/2^-$ isomers in $N = 82$ chain becomes different and highlights the dominance of $h_{11/2}$ orbital only (pure-seniority scheme). These isomers have also been compared with the identical trends for the $v = 2$, 12^+ isomers, and the $v = 3$, $33/2^+$ isomers in the $Z = 82$ chain. The identical behavior of all the high-spin isomers in various semi-magic chains strongly supports the goodness of seniority and generalized seniority up to very high-spin. This simple scheme of calculating the $B(EL)$ values may also be used to estimate the half-lives in unknown cases and, hence predict new isomers.

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准幻核中的近似广义辛弱数守恒

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摘要: 对称性在了解诸如原子核的转动、自旋和宇称、及同位旋等核结构性质中都起着重要的作用, 并且使复杂的原子核结构问题得以简化。辛弱数就是由于原子核的对相互作用中的对称性所导出的众所周知的好量子数。通过对丰中子和缺中子核素及核素的高自旋态的衰变数据分析来揭示辛弱数的近似守恒性质。研究表明, 在准幻核的高自旋同质异能素链中, 无论所涉及的价空间的核子轨道有何不同, 广义辛弱数总是近似的好量子数。

关键词: 广义辛弱数; 同质异能素; 准幻核

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